

Crossover from the pair contact process with diffusion to directed percolation

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Crossover behaviors from the pair contact process with diffusion (PCPD) and the driven PCPD (DPCPD) to the directed percolation (DP) are studied in one dimension by introducing a single particle annihilation and/or branching dynamics. The crossover exponents ϕ are estimated numerically as $1/\phi \approx 0.58 \pm 0.03$ for the PCPD and $1/\phi \approx 0.49 \pm 0.02$ for the DPCPD. Nontriviality of the PCPD crossover exponent strongly supports the non-DP nature of the PCPD critical scaling, which is further evidenced by the anomalous critical amplitude scaling near the PCPD point. In addition, we find that the DPCPD crossover is consistent with the mean field prediction of the tricritical DP class as expected.

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The absorbing phase transition (APT) has emerged during the last few decades as a prototype of nonequilibrium critical phenomena. The APT is a transition from an active phase into an inactive (absorbing) phase which is composed of absorbing (trapped) configurational states where the system cannot leave by the prescribed dynamic rules. As in the equilibrium critical phenomena, a few scaling exponents characterize and classify critical behaviors into universality classes [1,2]. Robustness of the directed percolation (DP) universality class with respect to the microscopic details led to the “DP conjecture” [3,4] that a model should belong to the DP class if it has a unique absorbing state without additional symmetry, conservation laws, quenched disorder, and long-range interactions. Although significant progress has been achieved to date, the full understanding of the main features of the APT universality classes is still far from complete.

To achieve such a goal, it is crucial to identify the universality class of the pair contact process with diffusion (PCPD) [4–6], which has been the most controversial topic these days, to our knowledge. The PCPD is a variant of the pair contact process [7] by allowing single-particle diffusion in addition to binary fission and annihilation dynamics ($2A \rightarrow 3A$ and $2A \rightarrow \emptyset$). Even with combined efforts of extensive and highly equipped numerical and analytical studies, its one-dimensional version has defied a consensus as yet [8]. Rather, various possible scenarios have been suggested with the DP class in the center of this controversy. In one scenario, the PCPD should belong to a universality class distinct from the DP with a unique set of critical exponents [9–11] or continuously varying exponents due to the marginal perturbation [12]. The other claims that the PCPD will eventually flow to the DP fixed point after a huge crossover time [13–15]. More scenarios can be found in Ref. [8]. In high dimensions, it is unquestionably clear that the PCPD and the DP exhibit different critical scaling.

The strong corrections to scaling are the main obstacle for the numerical study to determine the universality class of the PCPD. To make matters worse, the field theory with a single component order parameter is shown to be not viable [16]. It implies that the proper field theory, if it exists, needs (at least) two independent order parameters, which is also independently proposed by us from the numerical study of the driven PCPD (DPCPD) [10].

In this Rapid Communication, we propose an approach to settle down or at least moderate the PCPD controversy by scrutinizing closely the crossover behavior from the PCPD to the DP class. We allow unary branching and/or annihilation dynamics ($A \rightarrow 2A$ and $A \rightarrow \emptyset$) in the PCPD model, which should lead to the well-known DP critical dynamics as soon as it is introduced. The *nontrivial* nature of the DP critical line approaching the PCPD point as well as *anomalous* crossover scaling behavior would serve as an indicative of different scaling at the PCPD point, distinct from the DP.

Although the accurate estimation of the critical exponent values for the PCPD is severely plagued by corrections to scaling, the critical point itself can be located with relatively high precision. In Fig. 1, the density decay near criticality is shown for the PCPD model of which detailed dynamics is given in Eq. (1) with $w=0$. The system size at a tuning parameter $p=0.133\ 516$ ($0.133\ 522$) is 2^{18} (2^{20}) and the number of independent samples is 32 (104). All samples carry at least one neighboring pair until the end of simulations, $t = 10^9$ (10^8), which guarantees that finite size effects have not emerged as yet. From Fig. 1, the critical point is estimated as $p_c=0.133\ 519(3)$ with the number in the parentheses as the

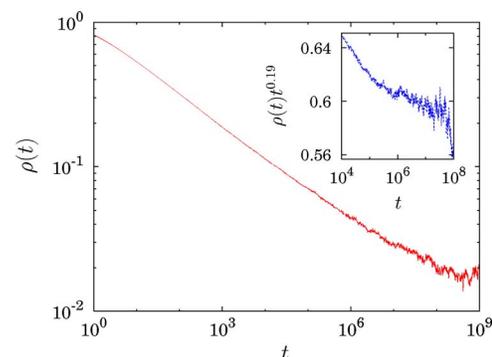
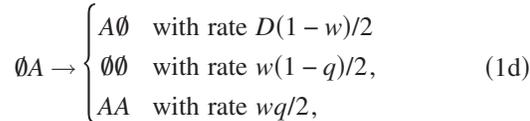
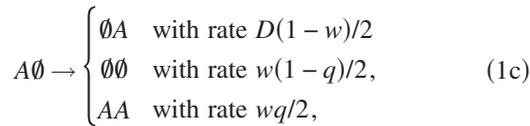
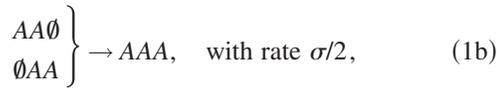
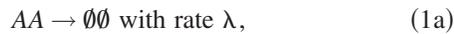


FIG. 1. (Color online) Log-log plot of the particle density decay for the PCPD at $p=0.133\ 516$. After $t \approx 10^8$, the density saturates, which signals that the system is in the active phase. Inset: Semi-logarithmic plot of $\rho(t)t^\delta$ versus t at $p=0.133\ 522$ with $\delta=0.19$. After $t \approx 10^7$, the curve veers down, which signals that the system is in the absorbing phase. This result is insensitive to the change of the δ value by up to 15%.

uncertainty of the last digit. It should be emphasized that our estimation of the critical point for the PCPD does not resort to any prescribed critical exponent value.

In the generalized crossover model including unary branching and/or annihilation processes, it is relatively easier to locate the critical points accurately by utilizing the known DP critical exponent values. With the annihilation $A \rightarrow \emptyset$ and/or the branching $A \rightarrow 2A$, the density decay in the absorbing phase should be exponential in time at least in one dimension, which nullifies the long-term memory effect of the isolated particles [12] and the simple unplugged DP scaling should be anticipated. The accurate information on the critical line as a function of the unary process rate w allows us to estimate the crossover exponent ϕ , which describes the crossover scaling, if it exists, from the PCPD to the DP.

The detailed evolution rules for the generalized crossover model are summarized using stoichiometric notations as



where $A(\emptyset)$ stands for a hard core particle (vacancy) and $0 \leq q \leq 1$. The periodic boundary conditions are employed on a one-dimensional lattice with size L . The case of $w=0$, $D=1$, $\lambda=p$, and $\sigma=1-p$ with a tuning parameter p corresponds to the PCPD model studied before [10], of which the critical point is accurately located through Fig. 1. In the case of nonzero w , the critical point is located by observing the flatness of $\rho(t)t^{\delta_{\text{DP}}}$ as a function of t with the DP critical exponent $\delta_{\text{DP}} \approx 0.1595$ [17]. In the following, we always set $D=1$ and $\lambda=1-\sigma=p$.

The simulation algorithm to mimic the dynamics in Eq. (1) is as follows: First, choose a particle at random. Then, choose one of its nearest neighboring sites randomly as a target site. If the target site is vacant, the chosen particle is annihilated with probability $w(1-q)$, branches a particle at the target site with probability wq , or hops to that site with probability $1-w$. The random selection of one of nearest neighbors amounts to the factor $\frac{1}{2}$ of the transition rates in Eqs. (1c) and (1d). In the case where the target site is already occupied, both particles are annihilated with probability p or one extra particle is generated at a randomly chosen nearest neighbor site of the pair with probability $1-p$ when that site is vacant. If the chosen site is already occupied, this branching attempt is rejected. After this update, the time increases by $1/N_t$, where N_t is the total number of particles at time t just before the update.

By observing how $\rho(t)t^{\delta_{\text{DP}}}$ behaves in the asymptotic re-

TABLE I. Critical point values $p_{c0}(w)$ and $p_{c1}(w)$ for various w 's for $q=0$ and $q=1$, respectively. The numbers in the parentheses indicate the uncertainty of the last digits.

| w | $p_{c0}(w)$ | $p_{c1}(w)$ |
|--------------------|--------------|--------------|
| 0 | 0.133 519(3) | 0.133 519(3) |
| 10^{-5} | 0.133 172(3) | 0.134 085(5) |
| 5×10^{-5} | 0.132 581(2) | |
| 10^{-4} | 0.132 080(3) | 0.135 81(1) |
| 2×10^{-4} | 0.131 325(5) | 0.136 98(1) |
| 3×10^{-4} | 0.130 715(5) | 0.137 92(1) |
| 4×10^{-4} | 0.130 185(5) | 0.138 73(1) |
| 5×10^{-4} | 0.129 708(3) | 0.139 46(1) |

gime, the critical point values p_c for nonzero w 's are estimated. If the system is in the active (absorbing) phase, those curves should veer up (down). Only at criticality, a flat straight line can be observed. Initially, all sites are occupied. The simulations were performed up to $t=10^7-10^8$ and no finite size effects have been observed with system size $L=2^{18}-2^{20}$. The numerical estimates for the critical point values $p_{c0}(w)$ at $q=0$ and $p_{c1}(w)$ at $q=1$ are tabulated in Table I. Figure 2 shows the critical density decay for various values of w at $q=0$ and 1. The flatness of the critical curves can be seen quite early even for very small w , while the PCPD lines ($w=0$) do not show any flatness within our observation time. It may be understood either by a tremendously long dynamic correction to the DP scaling [15] or by a non-DP scaling for the PCPD. In either case, it is evident that there exists a diverging crossover time scale as $w \rightarrow 0$.

The crossover scaling function near the PCPD critical point should take the form

$$\rho(w, \Delta; t) = t^{-\delta} F(\Delta w^{-1/\phi}, t w^{\mu_{\parallel}}), \quad (2)$$

where $\mu_{\parallel} = \nu_{\parallel}/\phi$ with the crossover exponent ϕ , ν_{\parallel} , and δ are the critical exponents of the PCPD, and $\Delta = p - p_c(0)$ with $p_c(0)$ being the PCPD critical point value. The crossover

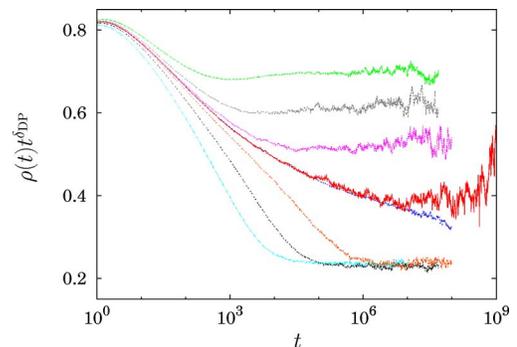


FIG. 2. (Color online) Semilogarithmic plots of $\rho(t)t^{\delta_{\text{DP}}}$ as a function of t for the generalized crossover model at criticality. The upper (lower) three curves correspond to $w=5 \times 10^{-4}$, 10^{-4} , and 10^{-5} from top (left) to bottom (right) at $q=0$ (1). The middle two curves are redrawn using the same data as in Fig. 1 for the PCPD ($w=0$) near the critical point.

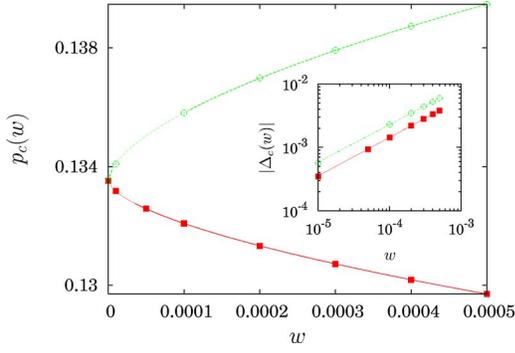


FIG. 3. (Color online) Critical points $p_c(w)$ for various w 's and their fitting functions. The lower (upper) curve corresponds to $q = 0$ (1). The error bar is smaller than the symbol size. The curves represent the fitting results of the critical points with the function $\Delta_c(w) = aw^{1/\phi}$. The estimated crossover exponent is $1/\phi = 0.58 \pm 0.03$. Inset: Log-log plot of $|\Delta_c(w)|$ vs w . The slope of the straight line is $1/\phi$.

time diverges as $w \rightarrow 0$ such that $\tau_{\text{cross}} \sim w^{-\mu_{\parallel}}$.

The crossover exponent ϕ can be calculated without the knowledge of the values of ν_{\parallel} and δ by studying how the DP critical line $p_c(w)$ approaches the PCPD point. This line should be one of the renormalization group flow lines such as $\Delta_c(w) \approx aw^{1/\phi}$ [18], where $\Delta_c(w) \equiv p_c(w) - p_c(0)$ and a is a (nonuniversal) constant. In Fig. 3, the crossover exponent is estimated using the least-square fitting, which turns out to have a nontrivial and universal value of $1/\phi = 0.58 \pm 0.03$. This nontrivial nature of the critical line ($\phi \neq 1$) signals strongly a possible non-DP scaling at the PCPD point.

The similar analysis can be applied to the DPCPD with diffusion bias by modifying the dynamics in Eqs. (1) in such a way that the dynamics (1d) is absent and the dynamics (1c) is replaced by

$$A\emptyset \rightarrow \begin{cases} \emptyset A & \text{with rate } (1-w) \\ \emptyset\emptyset & \text{with rate } w \end{cases}. \quad (1c')$$

With $w=0$, this model is the DPCPD studied in Ref. [10]. Since the diffusion bias should not change the DP scaling [10], this model with nonzero w is still expected to belong to the DP class. Figure 4 shows how the DPCPD crosses over to the DP as w increases and finds the crossover exponent $1/\phi \approx 0.49 \pm 0.02$, which is consistent with the mean field prediction of the tricritical DP class [19] and, in turn, confirms again that the upper critical dimension of the DPCPD is 1 [20].

To get a more concrete evidence for the non-DP nature of the PCPD scaling, we go back to the scaling function in Eq. (2), which can be rewritten at criticality as

$$t^{\delta_{\text{DP}}} \rho(w, \Delta_c; t) = t^{\delta_{\text{DP}} - \delta} F(a, w^{\mu_{\parallel}} t) = w^{\chi} \mathcal{G}(w^{\mu_{\parallel}} t), \quad (3)$$

where $\chi = \mu_{\parallel}(\delta - \delta_{\text{DP}})$ and $\mathcal{G}(x) = x^{\delta_{\text{DP}} - \delta} F(a, x)$. Since $\mathcal{G}(x)$ should be a finite constant as $x \rightarrow \infty$ to guarantee the DP scaling for nonzero w , the critical decay amplitude is proportional to w^{χ} for sufficiently small w where the crossover scaling is valid.

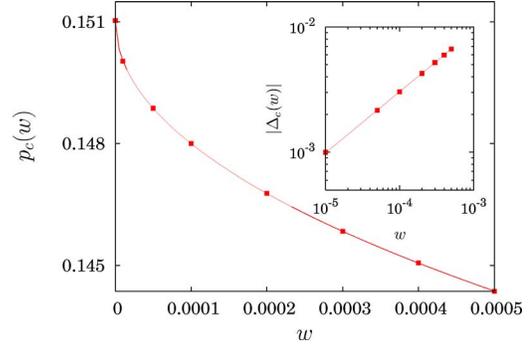


FIG. 4. (Color online) Phase boundary of the DPCPD crossover model with (1c'). Just as in Fig. 3, the crossover exponent is estimated by the least-square fitting as $1/\phi = 0.49 \pm 0.02$, which is consistent with the mean field crossover exponent [19]. Inset: Log-log plot of $|\Delta_c(w)|$ vs w .

If the PCPD belongs to the DP class ($\delta = \delta_{\text{DP}}$), the amplitude converges to a nonzero value as $w \rightarrow 0$. Otherwise, it decreases algebraically with $\chi \approx 0.03$, using the present best estimates for the PCPD exponent values as $\delta \approx 0.19$ and $\nu_{\parallel} \approx 1.85$ [9,10,15]. In Fig. 2, there are apparently two very different convergent behaviors for $q=0$ and 1. Naive estimates lead to $\chi \approx 0.08$ for $q=0$ and $\chi \approx 0.00$ for $q=1$, both of which do not simultaneously fit into either the DP or the non-DP scenario.

We find that this perplexing result originates from the narrowness of the crossover scaling region near the PCPD point. To get a hint for the valid scaling regime, we observe the pair density ρ_p as well as the particle density ρ in simulations. Both quantities should scale in the same way and their critical amplitudes also scale similarly as $\sim w^{\chi}$. Hence, the ratio $R(t) \equiv \rho(t)/\rho_p(t)$ should be independent of w for the valid crossover scaling regime in the asymptotic limit. The inset of Fig. 5 shows that both cases at $q=0$ and 1 have not

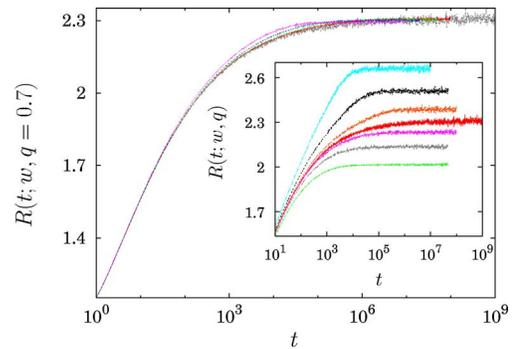


FIG. 5. (Color online) Density-pair density ratio $R(t)$ for various w 's with $q=0.7$ in the semilogarithmic scale. The critical points are located at $p_c = 0.134120(5)$, $0.133761(2)$, $0.133616(2)$, and $0.133563(2)$ for $w = 10^{-4}$, 2×10^{-5} , 4×10^{-6} , and 10^{-6} , respectively. The longest curve corresponds to the PCPD at $p = 0.133516$. The saturating values are almost identical for all w 's including the PCPD. Inset: Plots of $R(t)$ for w 's in Fig. 2. The upper (lower) three curves correspond to $q=1$ (0). For both cases, the ratio of two critical amplitudes approaches to the PCPD value as w becomes smaller.

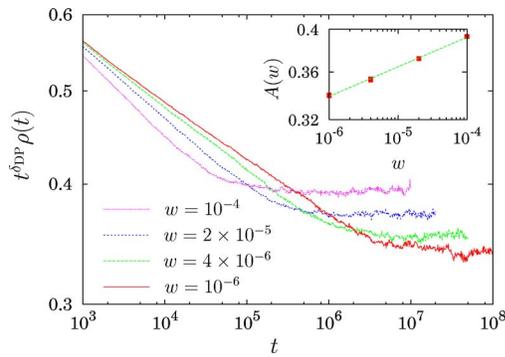


FIG. 6. (Color online) Log-log plots of $\rho(t)t^{\delta_{\text{DP}}}$ vs t for various w 's as in Fig. 5 with $q=0.7$ at criticality. Inset: Log-log plot of the critical amplitudes $A(w)$ vs w . The slope of the straight line is 0.03.

entered the scaling regime as yet up to $w=10^{-5}$, which implies that the naive estimates for χ should be significantly influenced by the corrections to scaling. Fortunately, we find a reasonably good scaling regime for not so small w at q

$=0.7$. Figure 5 shows that the scaling regime is reached already for $w=10^{-4}$. The inset of Fig. 6 shows how the critical amplitude of the density, say $A(w)$, behaves in the crossover scaling regime, from which we estimate $\chi=0.03(1)$ that is consistent with the PCPD estimate. This result strongly supports again the non-DP nature of the PCPD scaling. The crossover exponent estimated for the case of $q=0.7$ is $1/\phi \approx 0.57$, which is compatible with the previous estimation within error.

To conclude, we studied the crossover behavior from the pair contact process with diffusion and the driven pair contact process with diffusion to the directed percolation. We found that the crossover for the PCPD to the DP is nontrivial and the critical amplitude scaling is anomalous, which strongly supports the non-DP nature of the PCPD scaling. The crossover exponent for the DPCPD takes the mean field value of the tricritical DP (TDP) [19]. This implies that the two-dimensional PCPD might have connections to the TDP that is under our current investigation.

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